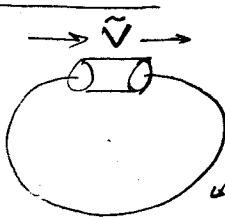


Longitudinal Motion

circular accelerator

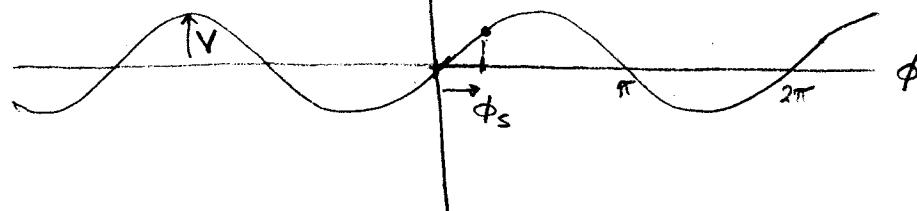


$$\tilde{V} = V \sin \omega_{RF} t$$

$$C = 2\pi R = n T_0 = n/f_0$$

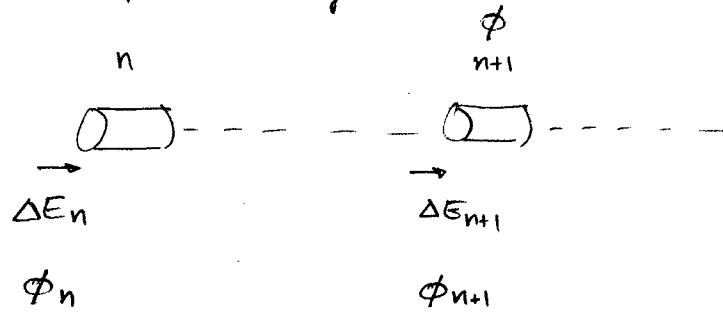
↑ assumes resonant cavity

let $\omega_{RF} T_0 = h$ $\omega_0 T_0 = 2\pi h$ $h = \text{harmonic \#}$

synchronous particle: E_s

$$\frac{dE_s}{dn} = eV \sin \phi_s$$

[note: must adjust beam field so that $\int B ds = \frac{P_s}{e} \cdot 2\pi$]

look at particle nearby: $E = E_s + \Delta E$ 

$$\text{easy to see } \Delta E_{n+1} = \Delta E_n + eV(\sin \phi_n - \sin \phi_s)$$

what about ϕ ?particle will take more/less time to go around, depending upon E (α, β)

$$T = \frac{C}{N}$$

$$\frac{\Delta T}{T} = \frac{\Delta C}{C} - \frac{\Delta N}{N}$$

$$\frac{\Delta N}{N} = \frac{1}{\gamma^2} \frac{\Delta P}{P}$$

$$\frac{\Delta P}{P} = \frac{1}{\beta^2} \frac{\Delta E}{E}$$

$$\Rightarrow \frac{\Delta T}{T} = \underbrace{\left(\frac{1}{\gamma_t^2} - \frac{1}{\gamma^2} \right)}_{\eta} \frac{\Delta P}{P}$$

($\eta = \text{slip factor}$)

$$\frac{\Delta C}{C} = \frac{1}{\gamma_t^2} \frac{\Delta P}{P} \text{ (more later)}$$

so, going once around, the phase will slip by an amount

$$\frac{\Delta\phi}{2\pi h} = \frac{\Delta T}{T} = n \frac{\Delta p}{p} = \frac{n}{\beta^2} \frac{\Delta E}{e}$$

$$\Rightarrow \phi_{n+1} = \phi_n + 2\pi h n \left(\frac{\Delta p}{p_s} \right)_{n+1}$$

or, we have

$$\Delta E_{n+1} = \Delta E_n + eV (\sin\phi_n - \sin\phi_s)$$

$$\phi'_{n+1} = \phi_n + \frac{2\pi h n}{\beta^2 E_s} \Delta E_{n+1}$$

(BASIC
program)

TRACK, or approximate w/ diff eq's ...

$$\frac{d\Delta E}{dr} = eV (\sin\phi - \sin\phi_s)$$

$$\frac{d\phi}{dn} = \frac{2\pi h n}{\beta^2 E_s} \Delta E$$

$$\Rightarrow \frac{d^2\phi}{dn^2} = \frac{2\pi h n eV}{\beta^2 E_s} (\sin\phi - \sin\phi_s)$$

Suppose $\phi = \phi_s + \Delta\phi$, $\Delta\phi \ll 1$, then

$$\begin{aligned} \sin\phi - \sin\phi_s &= \sin\phi_s \cos\Delta\phi + \sin\Delta\phi \cos\phi_s - \sin\phi_s \\ &\approx \Delta\phi \cos\phi_s \end{aligned}$$

$$\Rightarrow \frac{d^2\phi}{dn^2} - \frac{2\pi h n eV \cos\phi_s}{\beta^2 E_s} \Delta\phi = 0$$

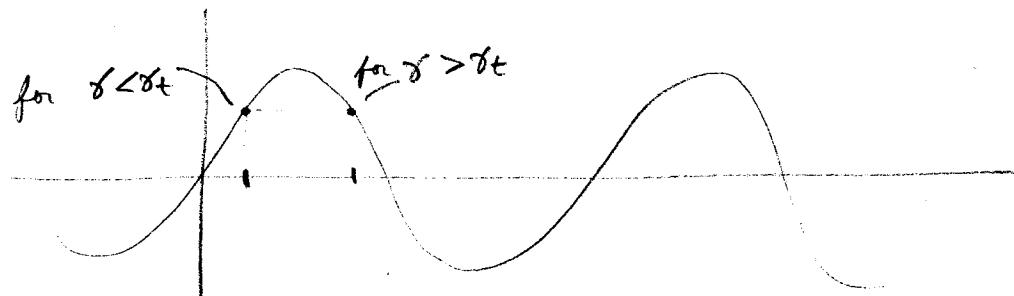
then motion is stable if

$$n \cos\phi_s < 0$$

(3)

So, when $\eta < 0$ [$\gamma < \gamma_t$], we want $\cos\phi_s > 0$

when $\eta > 0$ [$\gamma > \gamma_t$], we want $\cos\phi_s < 0$



\therefore if δ_t exists, need "phase jump @ transition"

$\gamma mc^2 = \text{transition energy}$

$$\text{BOO: } \gamma_t \approx 4.5$$

\therefore cross

$$\text{MI: } \gamma_t \approx 20$$

\therefore cross

$$\text{TEV: } \gamma_t \approx 18$$

\therefore do NOT cross

Small oscillation period...

$$\omega_s = \sqrt{-\frac{2\pi h \eta \text{ eV} \cos\phi_s}{\beta^2 E_s}}$$

$$\nu_s = \sqrt{-\frac{\hbar \eta \text{ eV} \cos\phi_s}{2\pi \beta^2 E_s}} = \text{synchrotron tune}$$

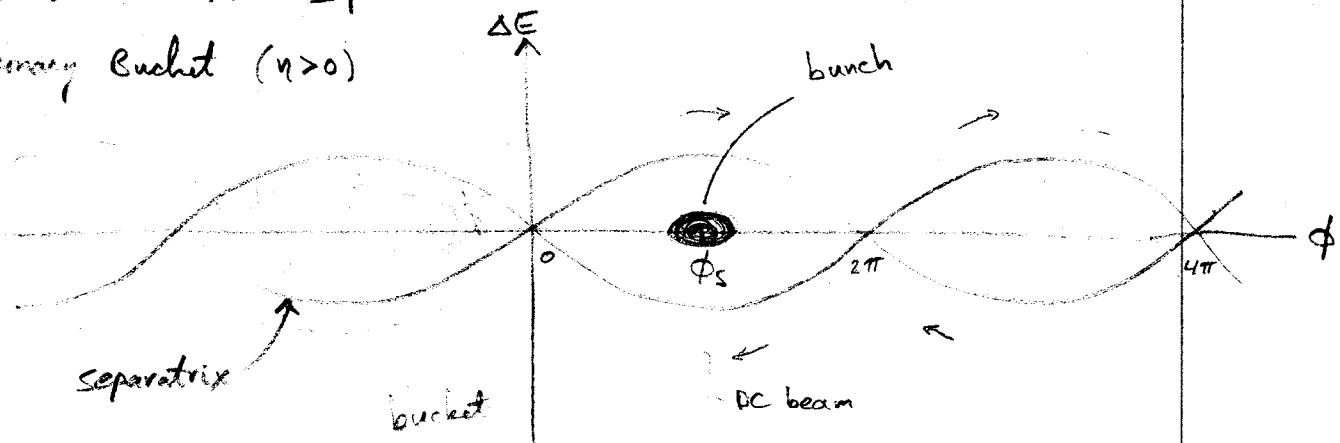
$$\frac{1}{\nu_s} = \text{synchrotron period (# turns)} \gg 1$$

ex: TEV @ 150 GeV

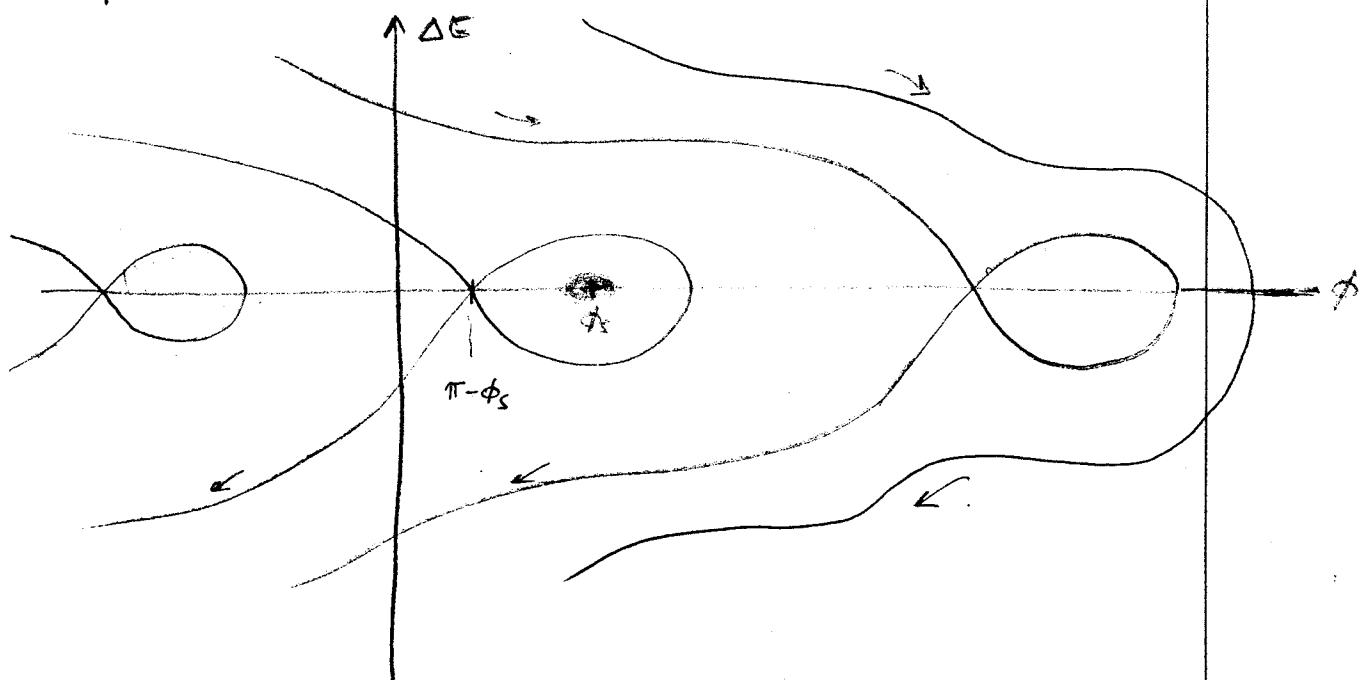
$$\nu_s = \sqrt{-\frac{(1113)(\frac{1}{18^2})(10^6 \text{ eV})(-1)}{2\pi \cdot 150 \cdot 10^9 \text{ eV}}} \approx \frac{1}{500}$$

$$\text{synchrotron "frequency"} = \nu_s f_0 = \frac{\nu_s c \beta}{C} \approx \frac{3 \cdot 10^8 \text{ m/s}}{500 \cdot 2\pi \cdot 10^3 \text{ m}} \approx 100 \text{ Hz}$$

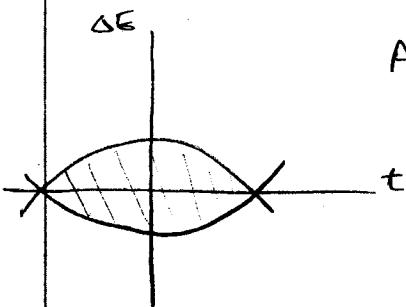
Longitudinal Phase Space

Stationary Bucket ($\eta > 0$)

v_s depends on amplitude of motion; near separatrix, $v_s \rightarrow 0$

Accelerating Bucket ($\eta > 0$)

Change to canonical variables: $\oint \mathbf{E} \cdot d\mathbf{t} = \text{adiabatic invariant}$

 ΔE 

Area of Stationary bucket:

$$A = \frac{8C}{\pi h c} \sqrt{\frac{eV \cdot G_s}{2\pi h |\ln|}}$$

(5)

Longitudinal Phase Space Area of beam = Longitudinal Emittance

ex: 95% longitudinal emittance of Gaussian Beam ...

assume: "small" oscillations, $\Delta\phi = \phi - \phi_s \ll 1$

then,

$$S = 2\pi^2 h \sqrt{-\frac{18\beta^2 E_s \text{ eV} \cos\phi_s}{\pi h \eta}} f_0 \sigma_t^2$$

$$= \frac{3}{h f_0} \sqrt{-\frac{2\pi h \eta E_s^3 \beta^2}{\text{eV} \cos\phi_s}} \left(\frac{\sigma_p}{p}\right)^2$$

ex: TEVATRON @ 150

$$A = \frac{8(2\pi \cdot 10^3 \text{ m})}{\pi(1113)(3 \cdot 10^8 \text{ m/s})} \sqrt{\frac{(10^6 \text{ eV})(150 \cdot 10^9 \text{ eV})}{2\pi(1113)\left(\frac{1}{18^2}\right)}} = 4 \text{ eV.sec}$$

$$S = 2\pi^2(1113) \sqrt{\frac{18(150 \cdot 10^9)(10^6)}{\pi(1113)\left(\frac{1}{18^2}\right)}} (47.7 \cdot 10^3) (2 \cdot 10^{-9})^2 \text{ eV.sec}$$

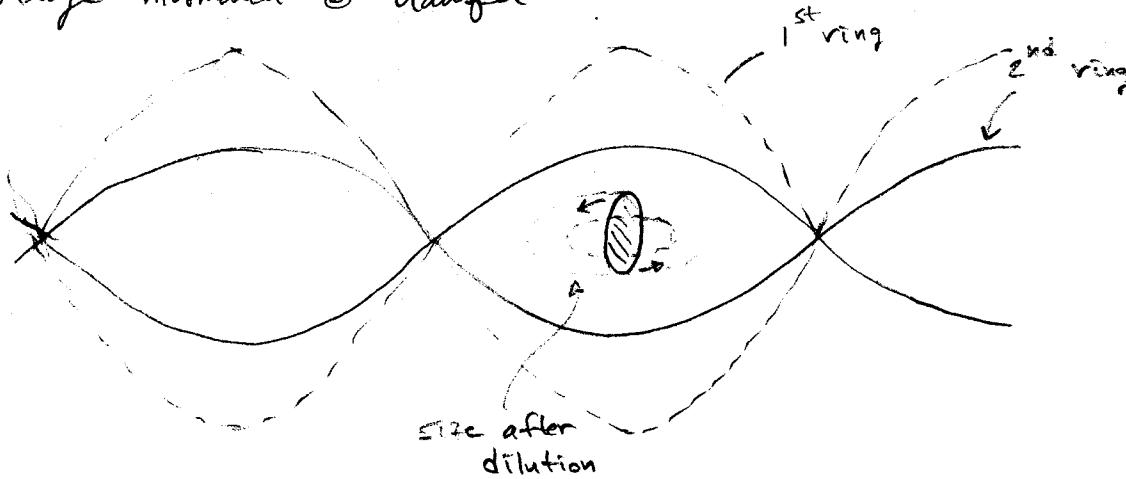
$$= 2 \text{ eV.sec}$$

note: $2 \text{ eV.sec} \Rightarrow \frac{\sigma_p}{p} = 0.36 \cdot 10^{-3}$

Machine matching: $\left(\frac{V}{h|\eta|}\right)_1 = \left(\frac{V}{h|\eta|}\right)_2$ on transfer

(assuming same f_{RF})

Voltage mismatch @ transfer

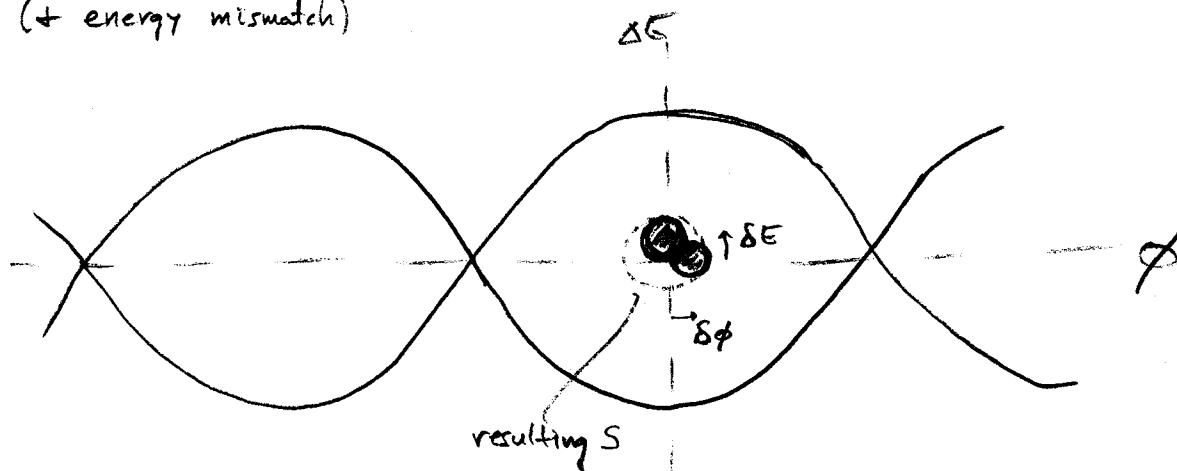


Increase in longitudinal emittance:

$$S/S_0 = \frac{1}{2} \cdot \frac{(1 + K/K_0)}{\sqrt{K/K_0}} \quad K = \frac{V}{h|\eta|}$$

Phase mismatch @ transfer is similar:

(+ energy mismatch)



$$S/S_0 = 1 + \frac{1}{2} \left(\frac{\delta E}{\sigma_E} \right)^2, \text{ or } = 1 + \frac{1}{2} \left(\frac{\delta \phi}{\sigma_\phi} \right)^2$$

[all these assume $S \ll A$]

Problems :

a) From $\frac{d\Delta E}{dn} = eV (\sin \phi - \sin \phi_s)$

$$\frac{d\phi}{dn} = \frac{2\pi h \eta}{\beta^2 E_s} \Delta E$$

Show that

$$\Delta E^2 + \frac{2\beta^2 E_s \text{eV}}{2\pi h \eta} (\cos \phi + \phi \sin \phi_s) = \text{constant}$$

- b) From a) above, and for $\phi_s = 180^\circ$, find an expression for the maximum ΔE on the separatrix (i.e., the "bucket height").

ANS: for a stationary bucket,

$$\Delta E_{\max} = \sqrt{\frac{2\beta^2 E_s \text{eV}}{\pi h \eta}}$$

- c) Evaluate the bucket height for the Tevatron at
- i) 150 GeV
 - ii) 980 GeV